

# Relaxation

Relaxation times:  $T_1$  and  $T_2$

Relaxation rate =  $1/T_{1,2}$



Relaxation in NMR

## Relaxation spin-lattice:

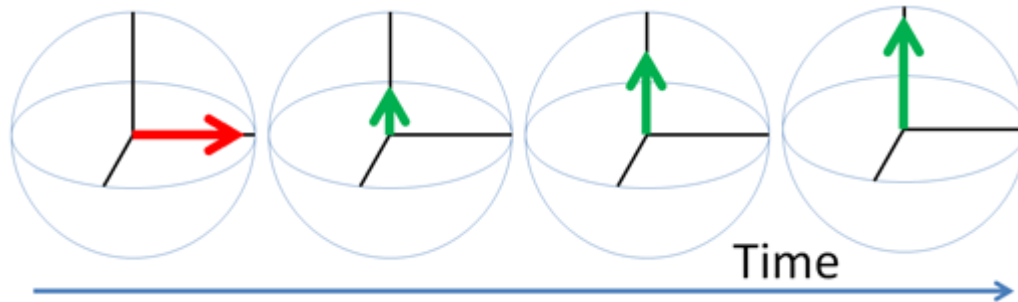
Transfert of energy from the system of spins to an acceptor of the lattice: rotations, translations, vibrations fluctuating in the domain of larmor frequencies

## Relaxation spin-spin:

Destruction of the component  $M_{xy}$  after a pulse.

## Two mechanisms: spin-lattice and spin-spin relaxation

**Spin-lattice relaxation** ( $T_1$ ) is the mechanism by which an excited magnetization vector returns to equilibrium along the axis of the static applied magnetic field.



The relaxation process is induced by field fluctuation due to molecular motion. (The local field experienced by a molecule changes when the molecule reorients)

**The correlation time** - $\tau_c$  (Tau-c): During this time the system does not move

The magnitude of the relaxation time depends highly on:

- the type of nuclei (nuclei with spin  $1/2$  and low magnetogyric ratio have usually long relaxation time whereas nuclei with spin  $>1/2$  have very short relaxation time)

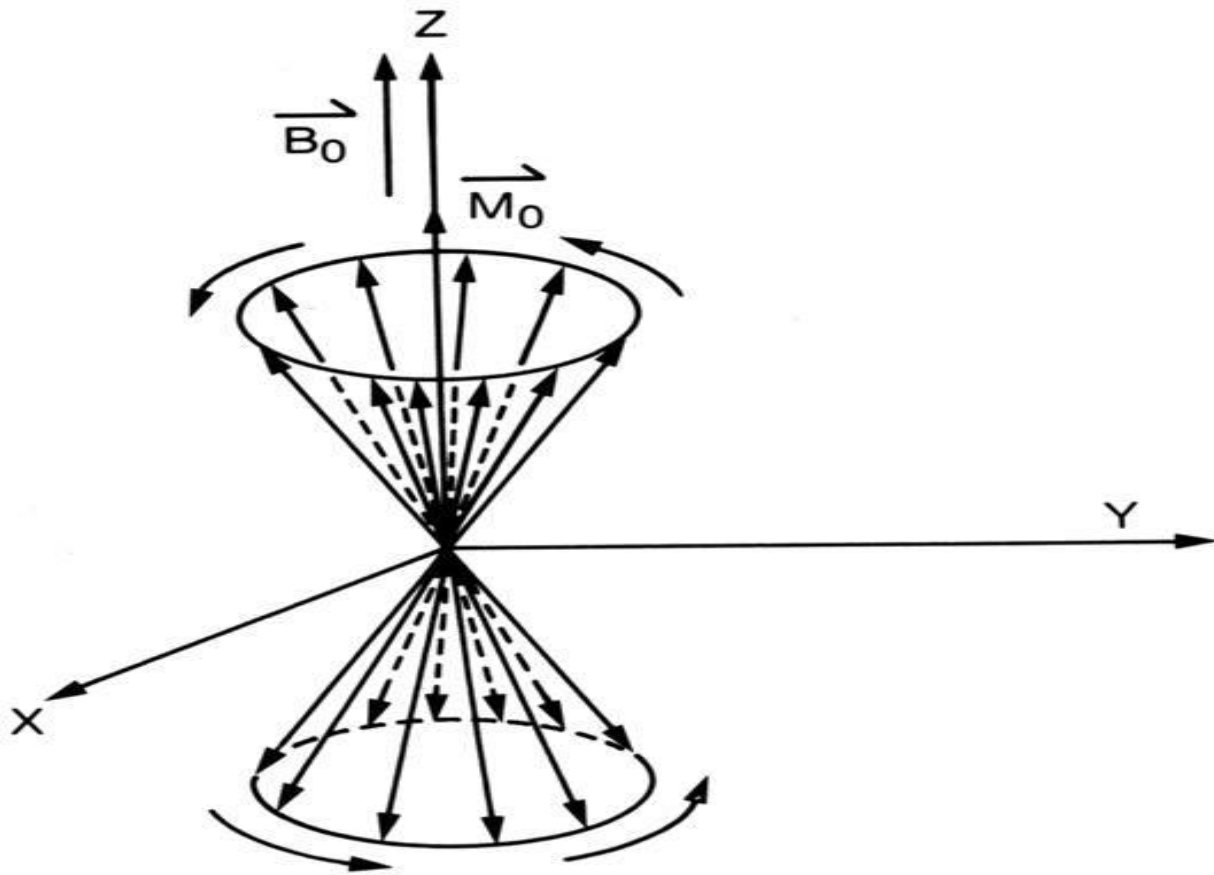
- other factors like the physical state (solid or liquid state), the viscosity of the solution, the temperature ... etc.

In other words the relaxation time depends on the motion of the molecules.

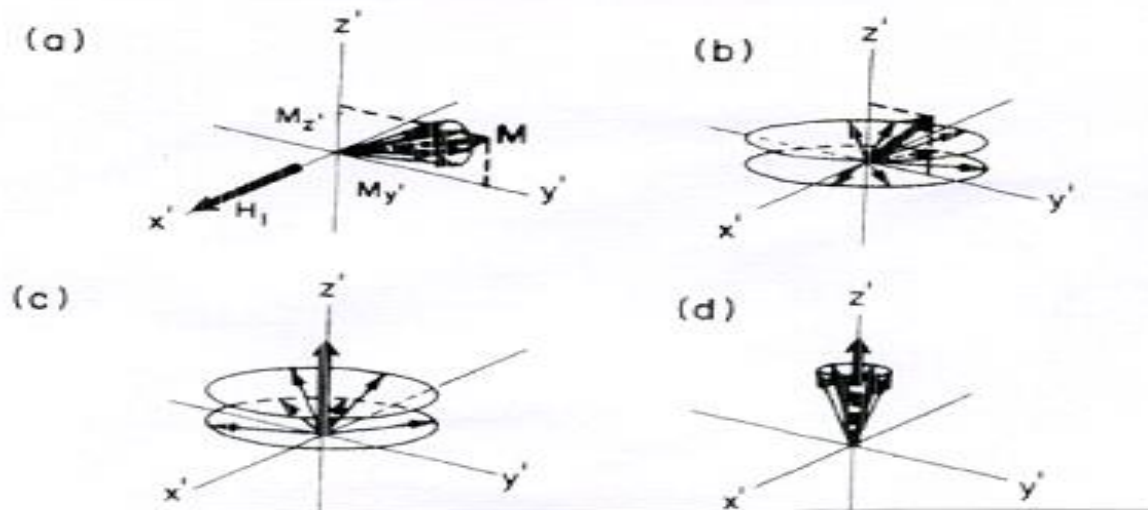
### Definitions:

The correlation time -  $\tau_c$  (Tau-c): represents the time it takes for a molecule to reorient by 1 degree ("tumbling time").

The spectral density -  $J(\omega)$ : describes the ranges of frequency motion that are present. Not all molecules tumble at a unique rate: molecules tumble, collide, change direction... at a range of rates up to the maximum rate of  $(1/\tau_c)$ .



## Pulsed NMR- Relaxation



The system evolves under two relaxation mechanisms:

- For a set of spins coupled by interactions, as in a solid, the system can tend very rapidly towards an internal quasi-equilibrium with a time constant,  **$T_2$ , named spin-spin relaxation time**. During  $T_2$  there is disappearance of the transversal magnetization by dispersion of the spins in the XOY plane.
- For such a set of spins in contact with the lattice, there is a **spin-lattice relaxation** characterized by the time  **$T_1$** , and then restoration of the magnetization  $M_0$  along  $B_0$ .

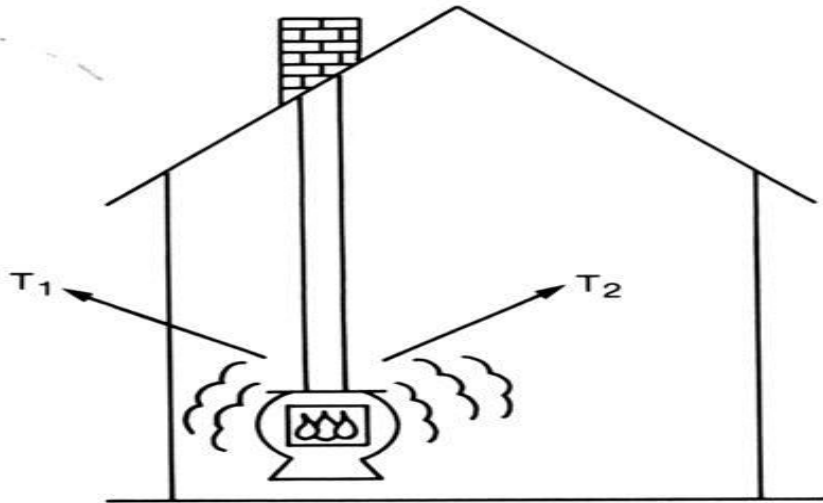
Generally, and specially for solids,  $T_2 \ll T_1$

Bloch's equations:  $dM_z / dt = - (M_z - M_0) / T_1$

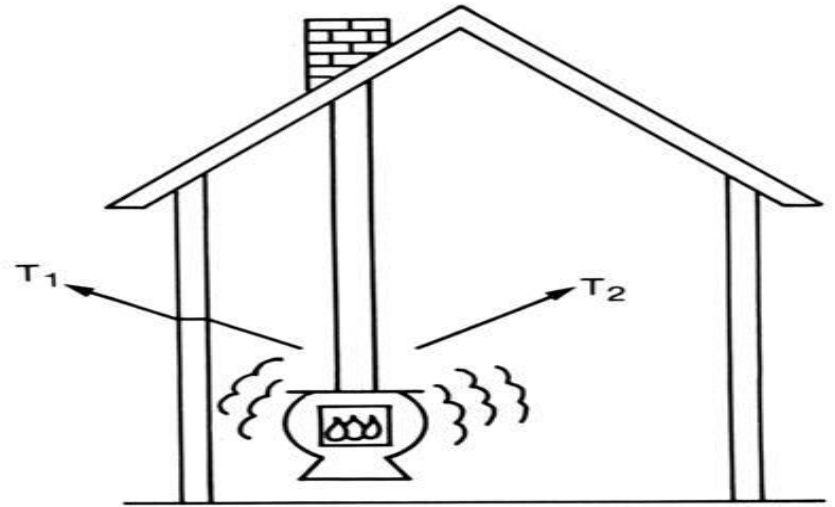
$dM_x / dt = -M_x / T_2$      $dM_y / dt = -M_y / T_2$

Stove-in-a-house analogy to T1 and T2 relaxation. In (a), the house is not insulated. So the energy emitted by the stove is lost to the surrounding environment (lattice) as fast as it is dispersed throughout the house. In (b), the house is insulated so that the energy from the stove is not dissipated to the surroundings nearly as fast as it equilibrates in the house.

(a)  $T_1^{-1} = T_2^{-1}$



(b)  $T_1^{-1} < T_2^{-1}$





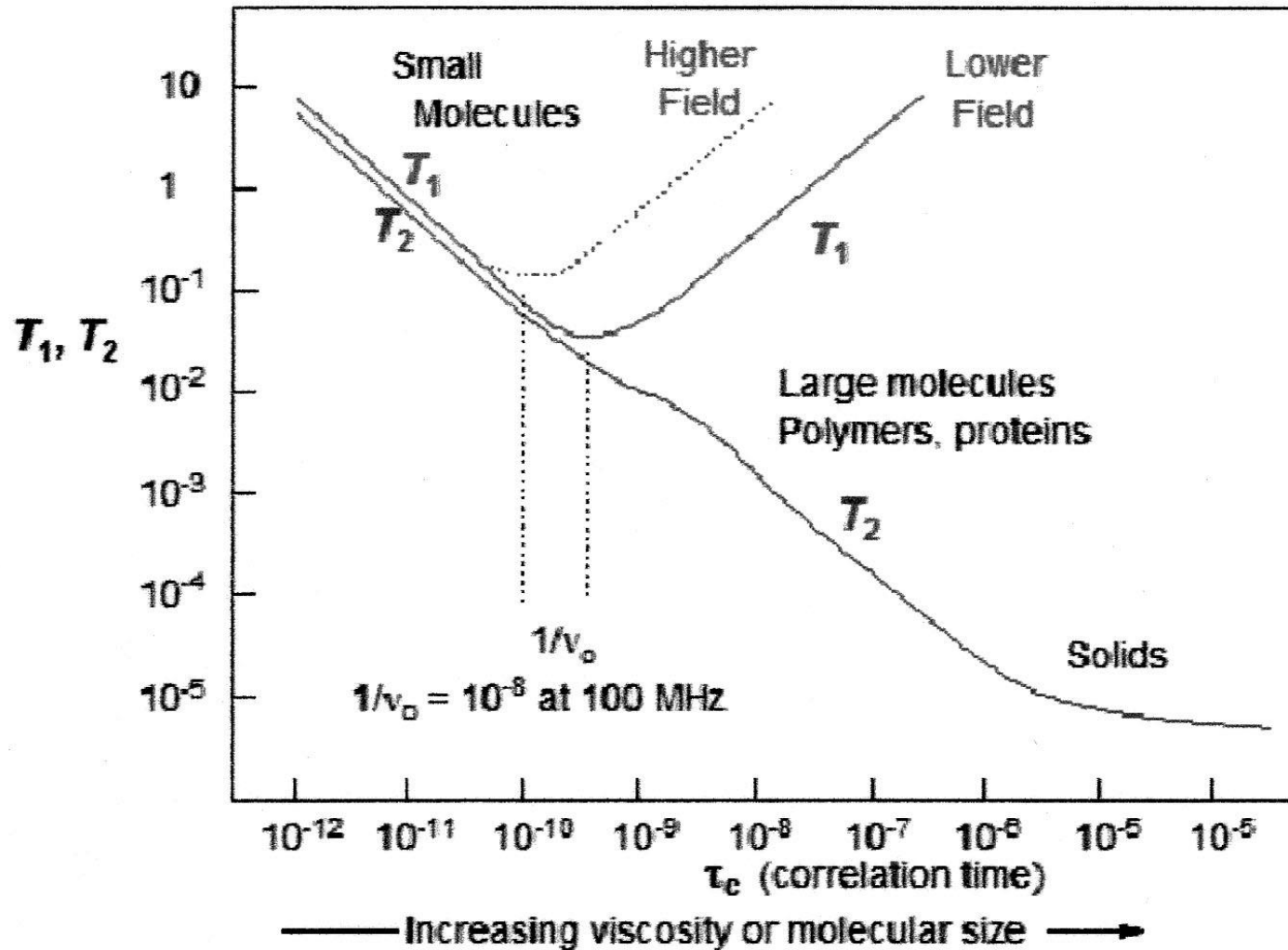
Interaction	Range of interaction (Hz)	Relevant parameters
<b>Dipolar coupling</b>	$10^4 - 10^5$	<ul style="list-style-type: none"> <li>- abundance of magnetically active nuclei</li> <li>- size of the magnetogyric ratio</li> </ul> $1/T_1 = R_1 = k \cdot \gamma_I^2 \cdot \gamma_S^2 \cdot (r_{IS})^2 \cdot f(\tau)$
<b>Quadrupolar coupling</b>	$10^6 - 10^9$	<ul style="list-style-type: none"> <li>- size of quadrupolar coupling constant</li> <li>- electric field gradient at the nucleus</li> </ul>
<b>Paramagnetic</b>	$10^7 - 10^8$	concentration of unpaired electrons (paramagnetic impurities)
<b>Scalar coupling</b>	$10 - 10^3$	size of the scalar coupling constant (interacting nuclei with close frequency)
<b>Chemical shift anisotropy (CSA)</b>	$10 - 10^4$	<ul style="list-style-type: none"> <li>- size of the chemical shift anisotropy</li> <li>- symmetry at the nuclear site</li> </ul>
<b>Spin rotation</b>		<p>Intramolecular dynamic process. Most effective for small, symmetrical molecules, or freely rotation methyl groups.</p> $1/T_1 = R_1 \uparrow \text{ with } T \uparrow$ <p>Coupling of the nuclear spin angular momentum with the molecular rotation angular momentum.</p>

## Bloch equations

$$\frac{dM_z}{dt} = \frac{M_0 - M_z}{T_1}$$

$$\frac{dM_{X(Y)}}{dt} = \frac{M_{X(Y)}}{T_2}$$

# Variation of $T_1$ and $T_2$ against viscosity or molecular size



## Measurement of $T_1$ with $180^\circ$ , $\tau$ , $90^\circ$ sequences

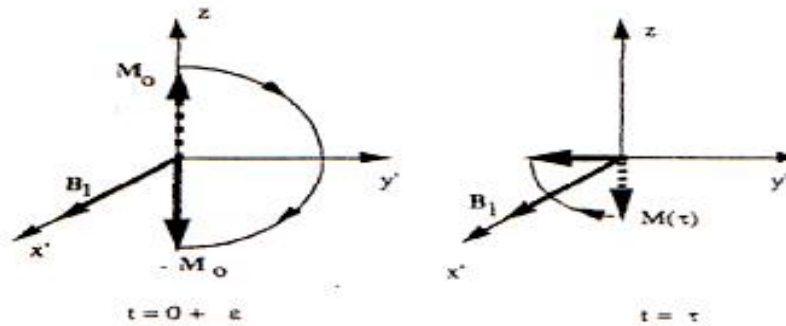


Figure (a)

Figure (b)

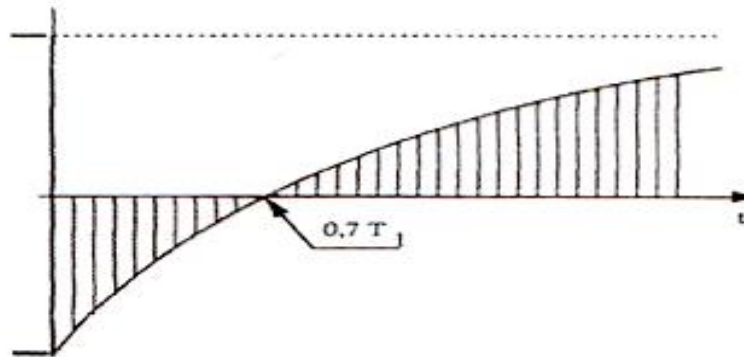
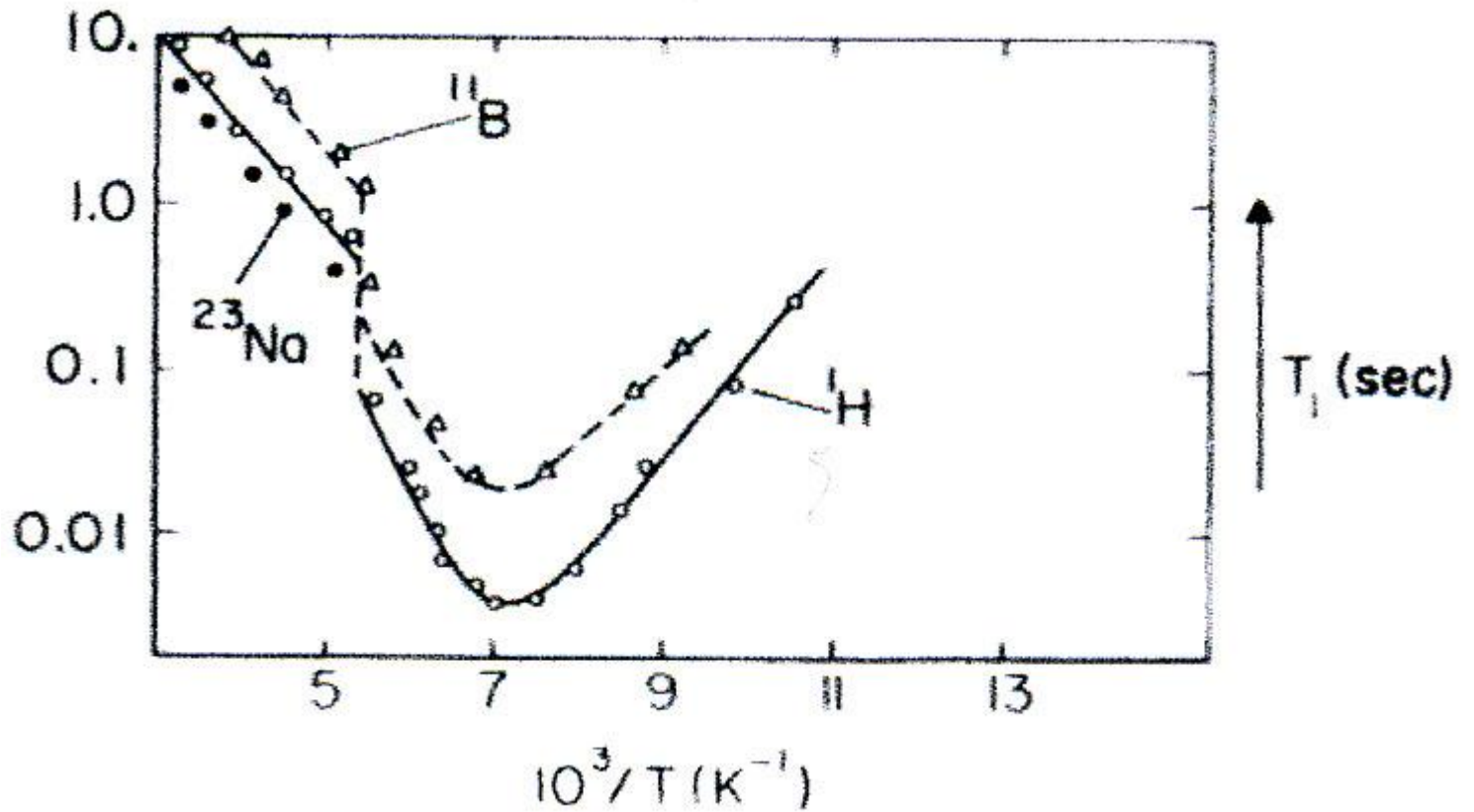


Figure (c)

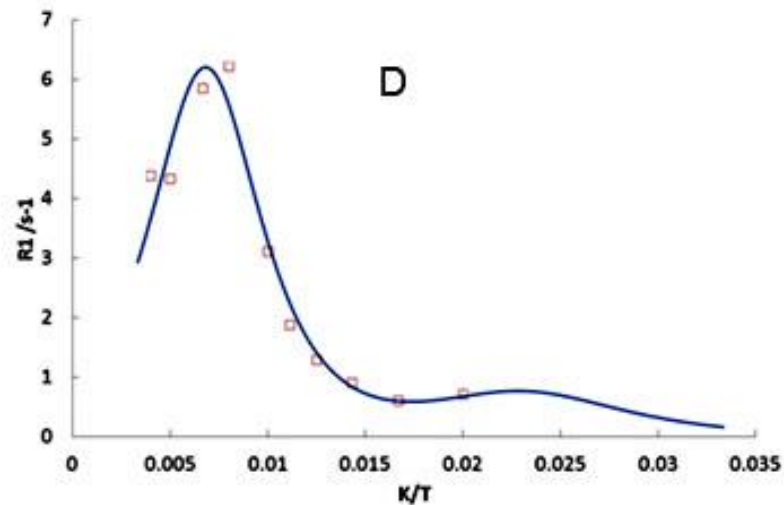
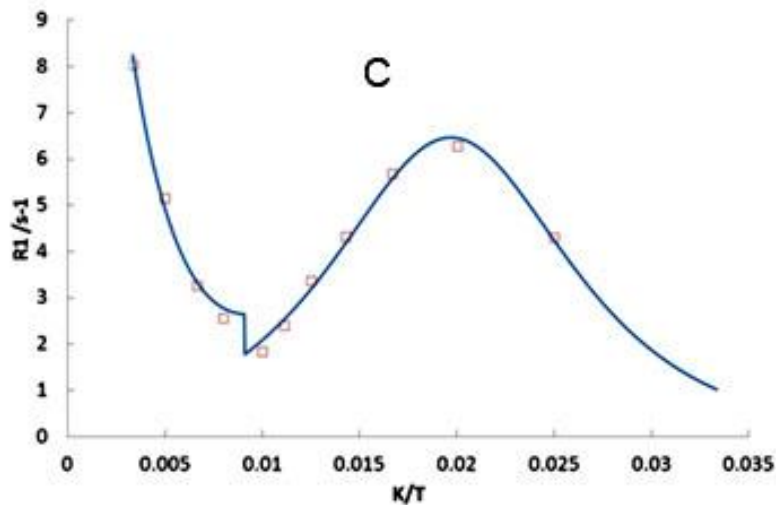
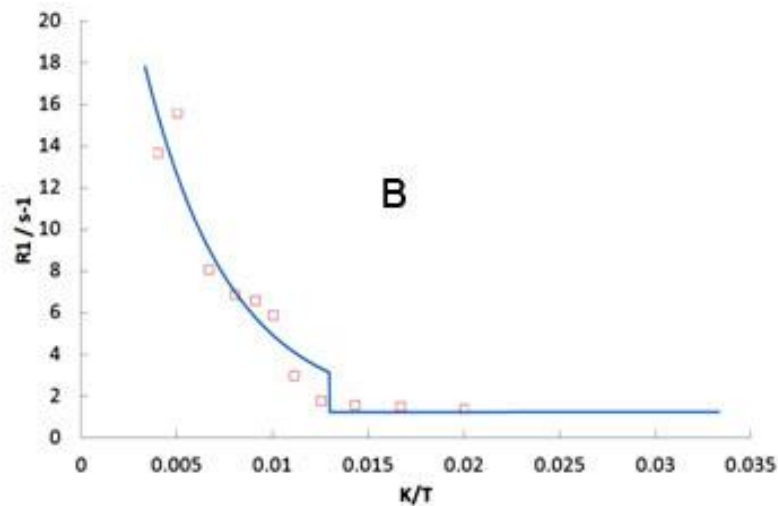
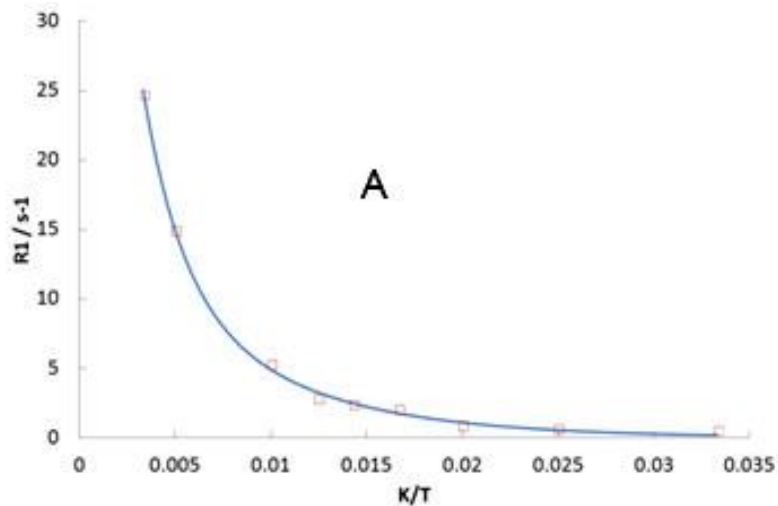
a)  $M$  is inverted by a  $180^\circ$  pulse at time 0; b) After a time  $\tau$  a  $90^\circ$  rotates  $M$  to the  $Y'$  (or  $-Y'$ ) axis. c) the amplitude of the FID after the  $90^\circ$  pulse, which is proportional to the amplitude of  $M$  at time  $\tau$ , is plotted as a function of  $\tau$ .

$$dM_z/dt = -(M_z - M_0)/T_1 \Rightarrow M_z = M_0 (1 - 2 \exp(-t/T_1))$$

Spin-lattice  $T_1$  in  $\text{NaBH}_4$  vs  $1/T$  showing a phase transition at about  $190^\circ\text{K}$



Relaxation rate  $R_1$  versus  $T^{-1}$ . A- SBA-80: spin-rotation;  
 B- Zeolite Omega: spin-rotation + dipolar contribution (cut-off 77K)  
 C- Zeolite ZSM5: spin-rotation + dipolar contribution (cut-off 110K) ;  
 D- Zeolite Rho: two dipolar contributions



# Relaxation due to dipolar interaction (spins $\frac{1}{2}$ )

$$\frac{1}{T_1} = K \left[ \frac{\tau_c}{1 + \omega_0^2 \tau_c^2} + \frac{4\tau_c}{1 + 4\omega_0^2 \tau_c^2} \right]$$

$$\frac{1}{T_2} = \frac{K}{2} \left[ 3\tau_c + \frac{5\tau_c}{1 + \omega_0^2 \tau_c^2} + \frac{2\tau_c}{1 + 4\omega_0^2 \tau_c^2} \right]$$

$$K = f(\tau_c, \gamma, \omega_0, r^{-6})$$

$\omega_0$  = Larmor frequency

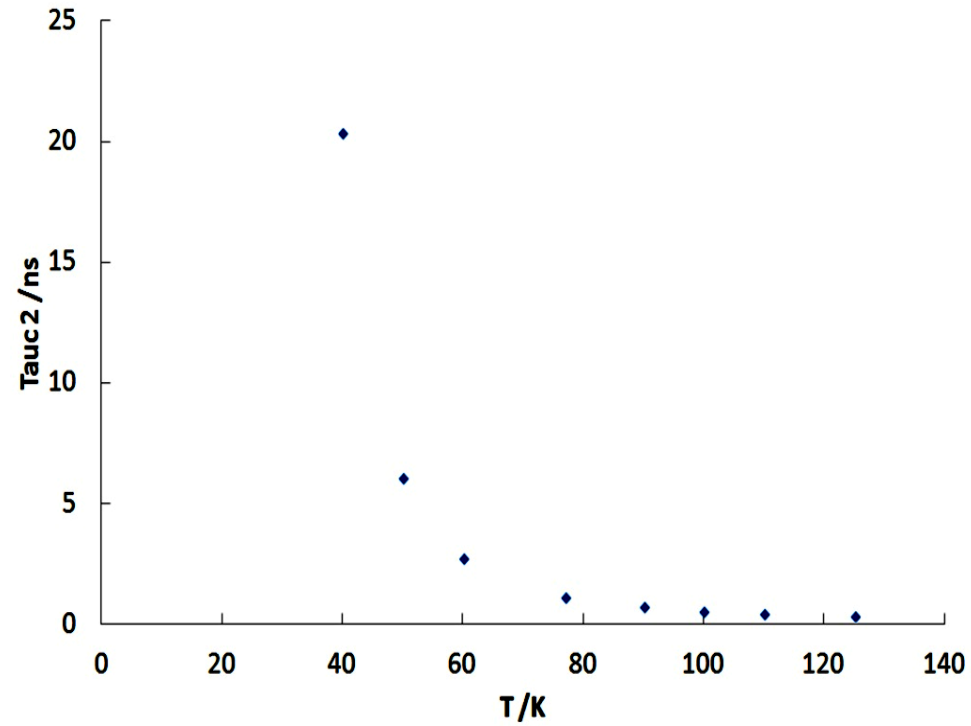
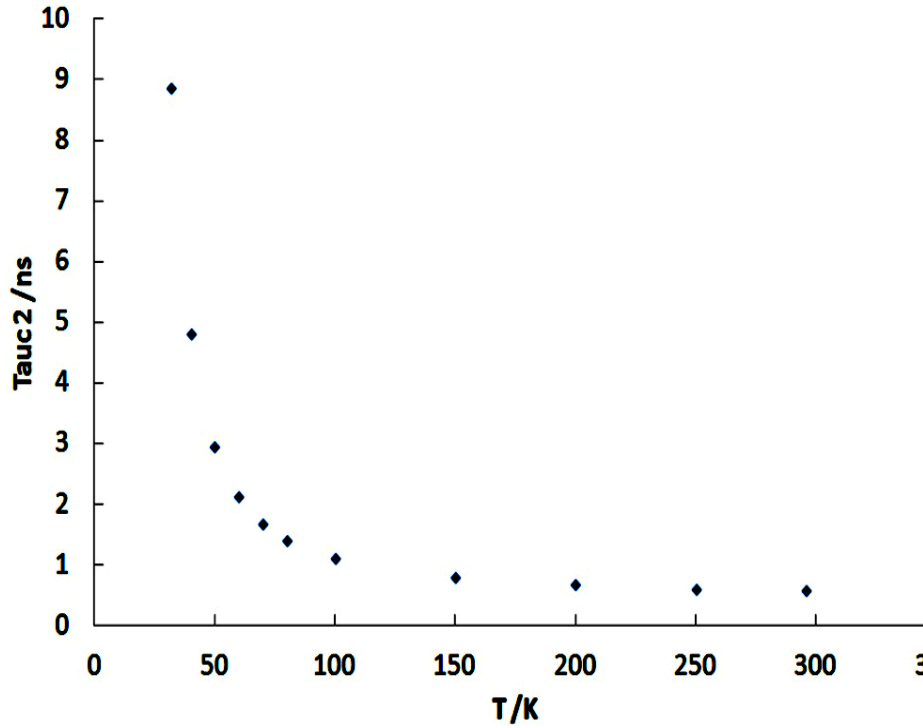
$\tau_c$  = correlation time of the tumbling motion

$\gamma$  = gyromagnetic ratio

# The dipolar correlation time of the N<sub>2</sub> molecule *versus* T within

Y

mordenite

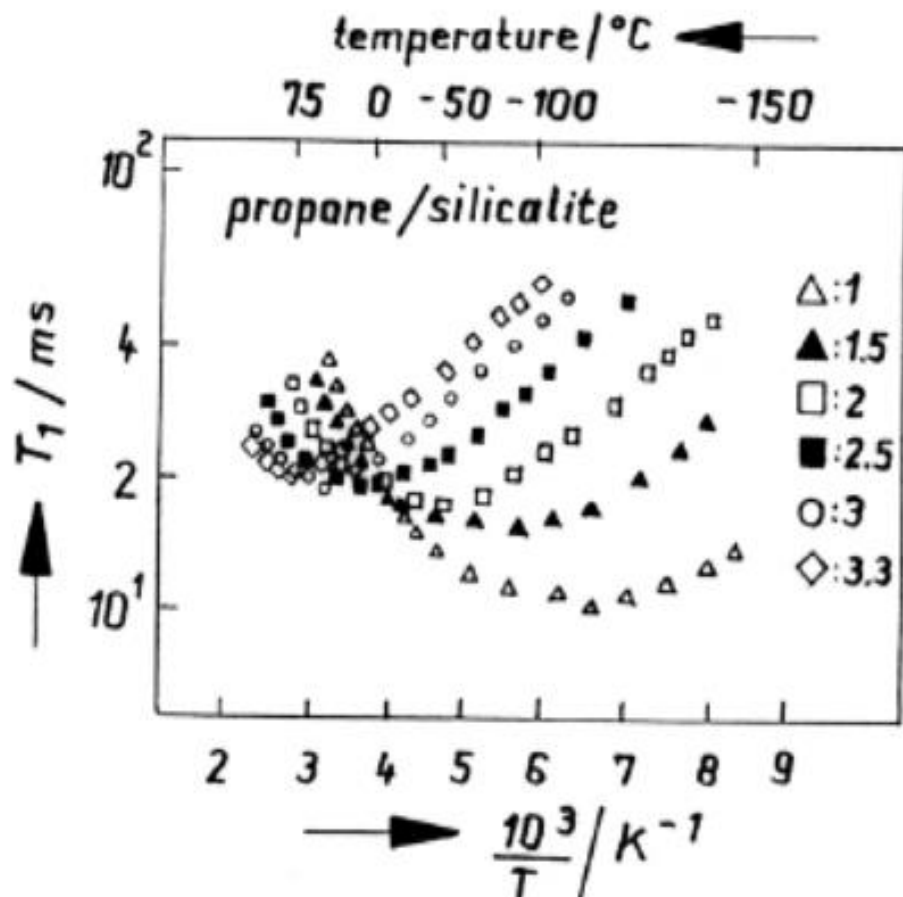
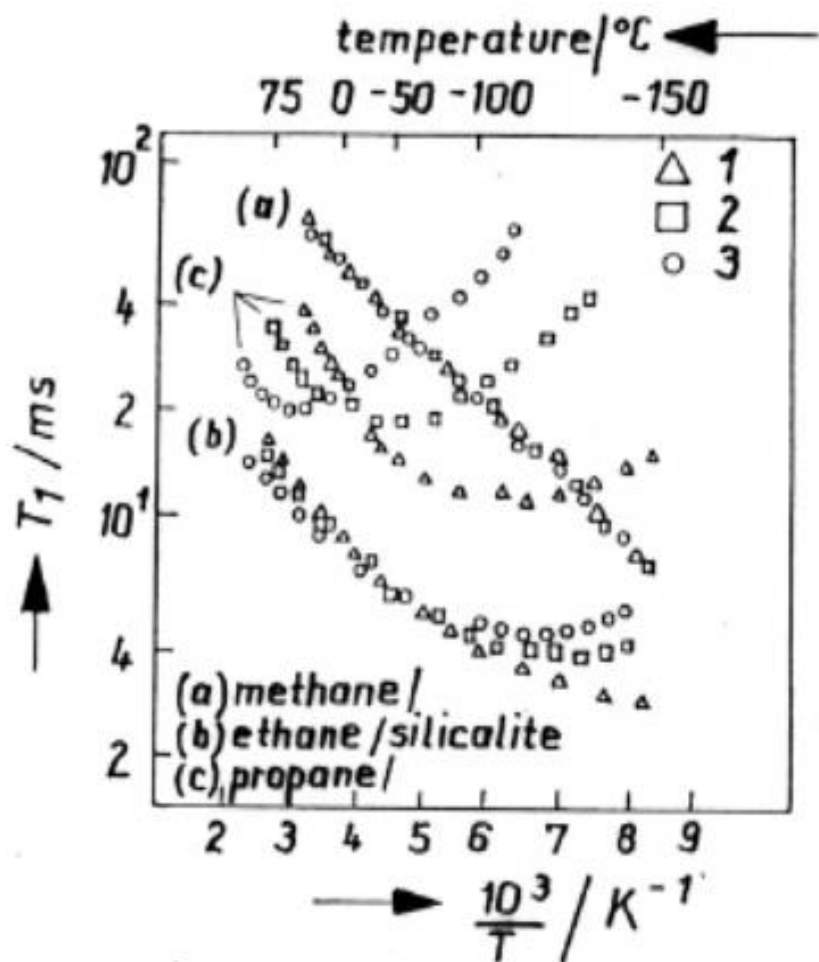




# Variation of T1 versus 1/T of adsorbed molecules

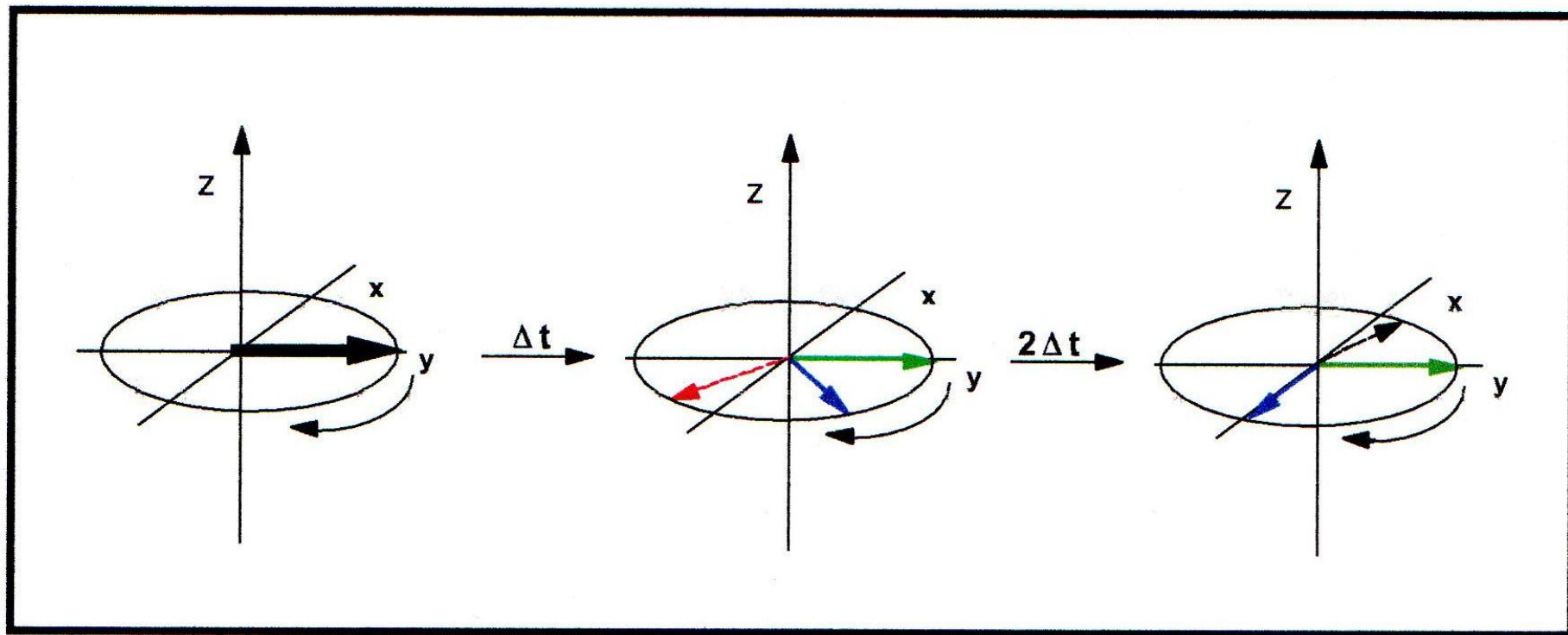
a: methane; b: ethane; c: propane

1  $\Delta$ , 2  $\square$ , 3  $\circ$  molecules / channel interaction



- Hahn Spin-echo

Spins processing at different velocities because of different chemical shifts are color coded



When a rabbit and a turtle run in one direction for a certain time, then turn around and run in the opposite direction with the same speed for the same time they will arrive at the starting point at the same time.

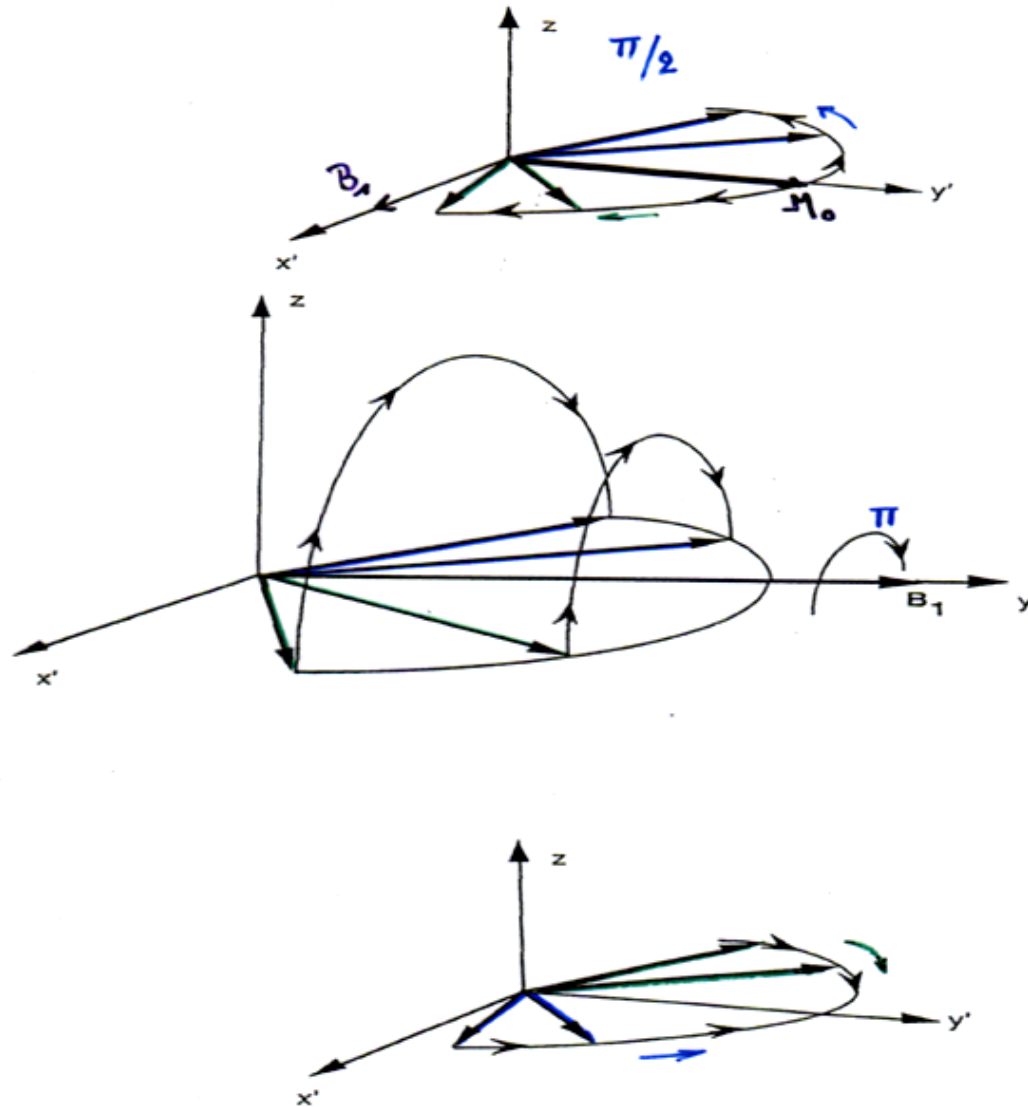




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# Hahn spin-echo experiment

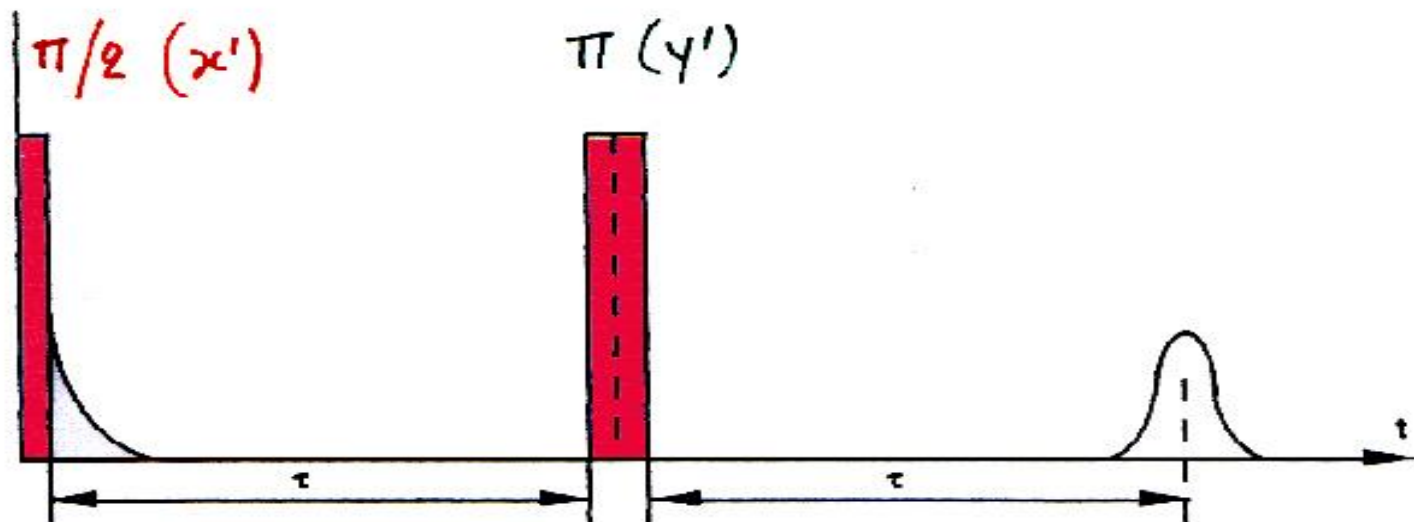
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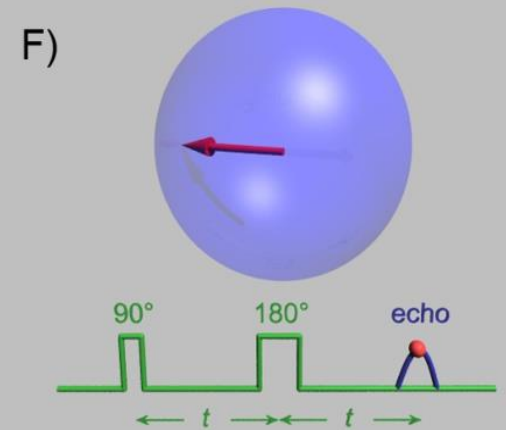
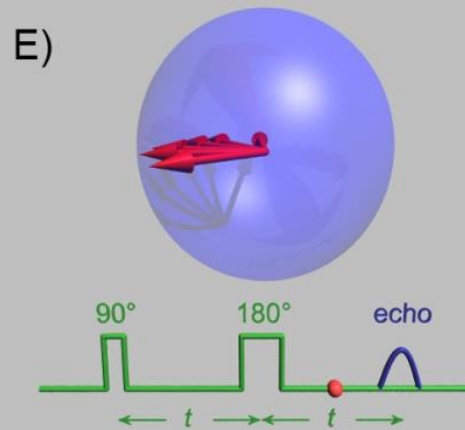
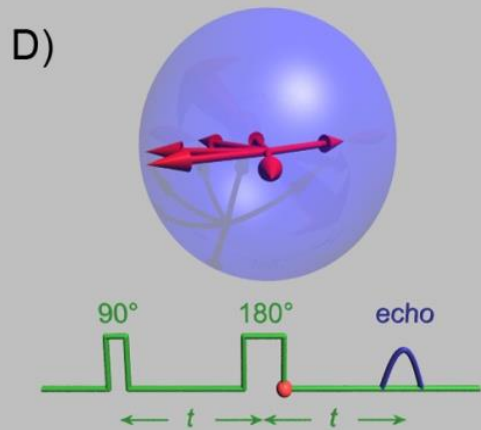
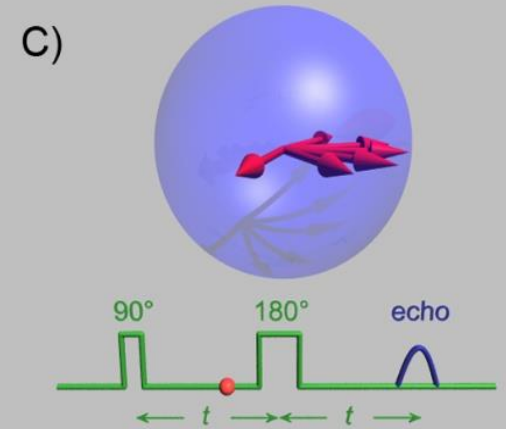
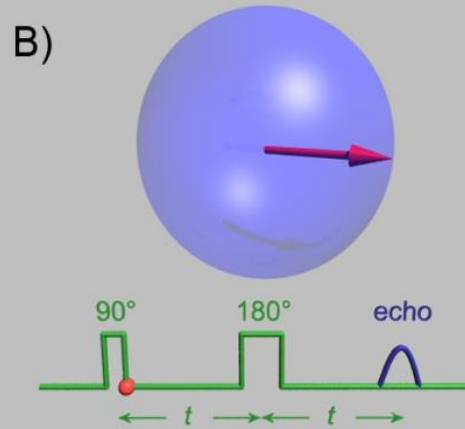
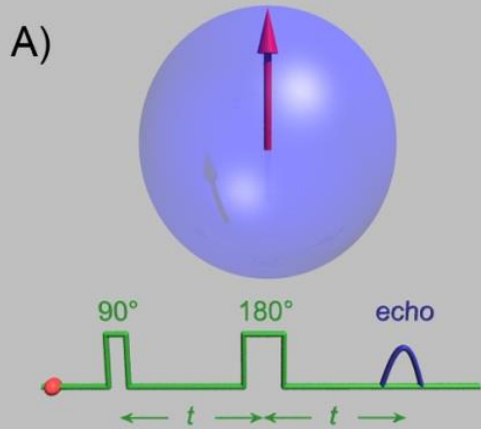
## Hahn spin-echo experiment

The distribution of the spin in the  $X'OY'$  plan after a pulse  $\pi/2$  is due to the spin-spin interactions but also to the inhomogeneity of  $B_0$ . The spin-echo experiment overcomes the inhomogeneity problem. In addition it is used for several applications of NMR.

$90^\circ, \tau, 180^\circ$  sequence, and observation at time  $2\tau$  of a free induction "echo".

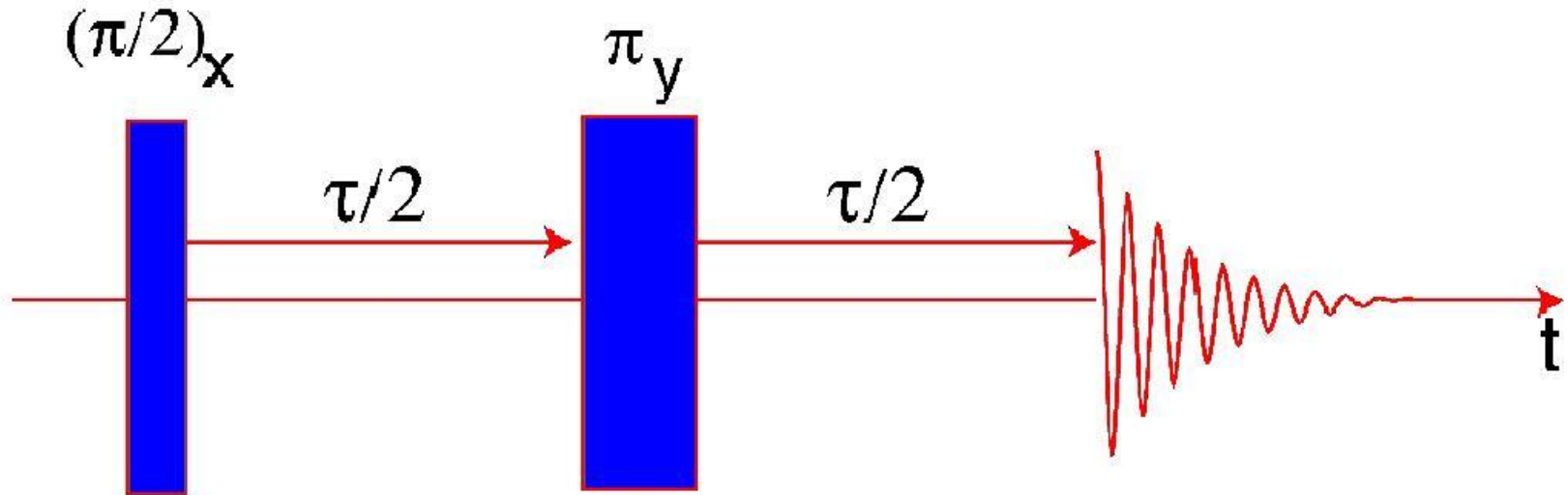


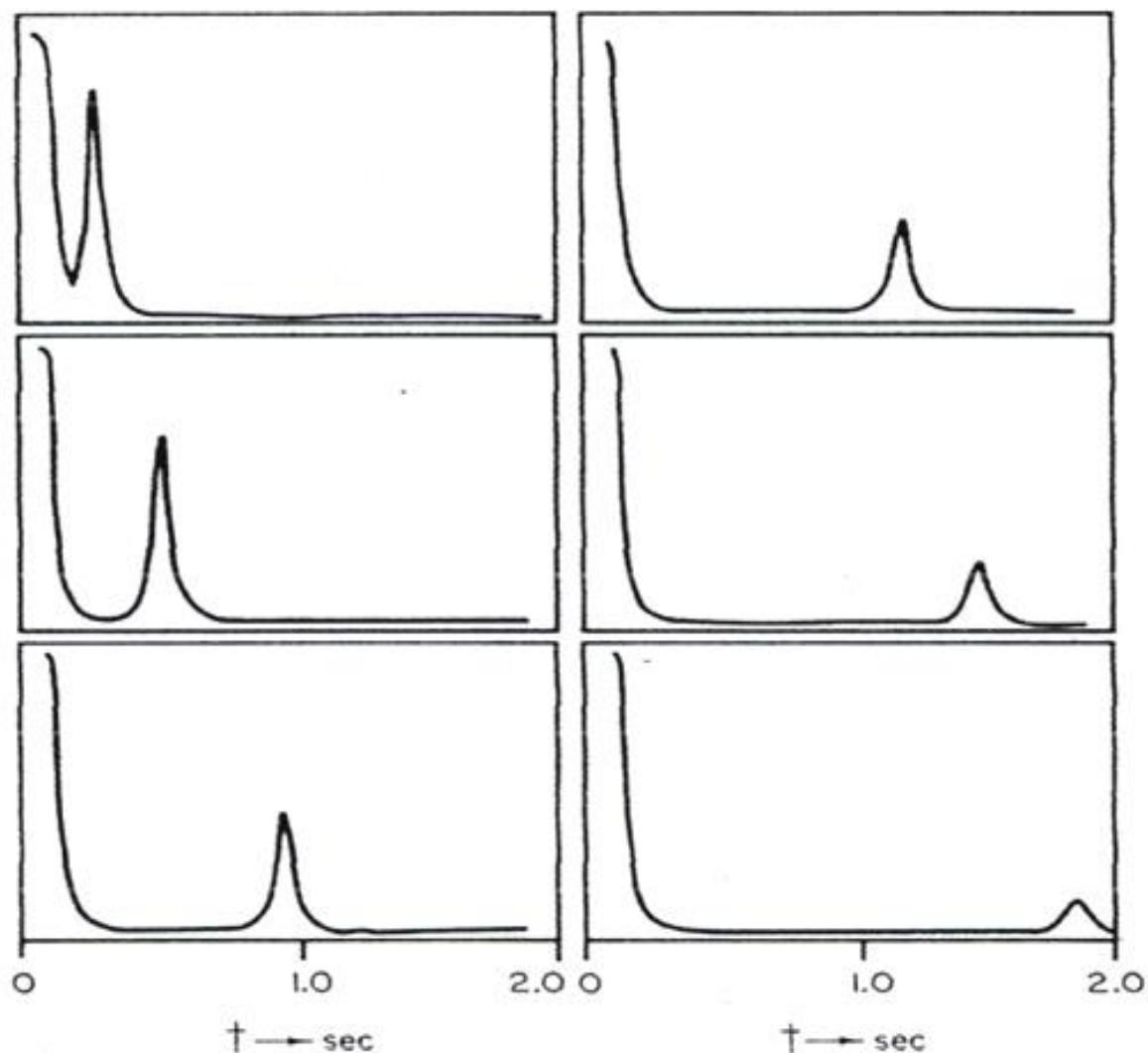
# Spin-echo experiment





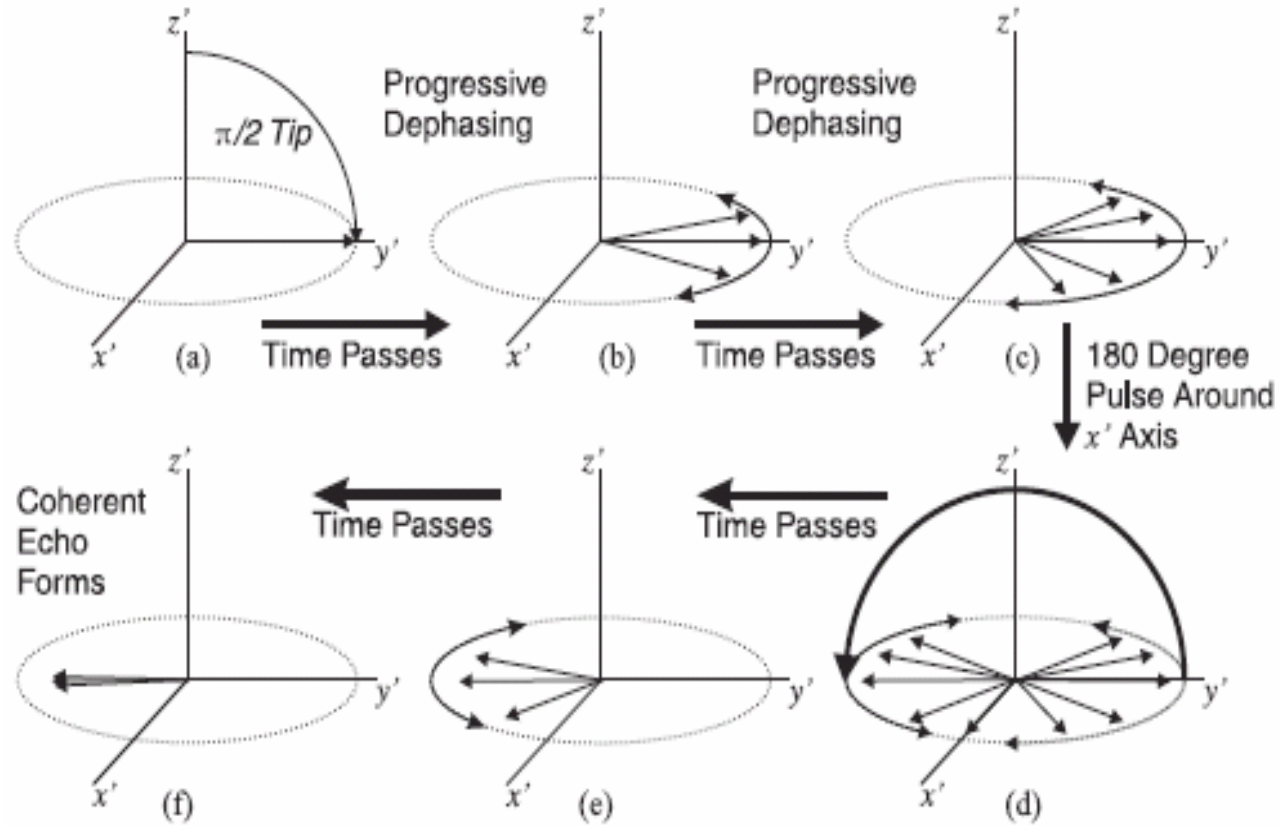
# Spin-echo experiment



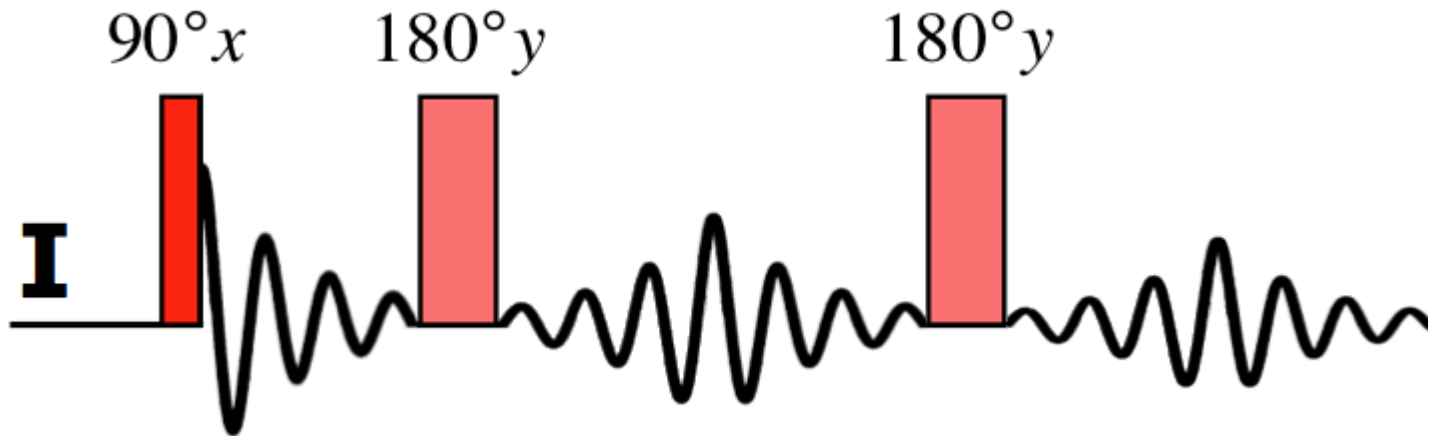


**Figure 2.5.** A typical Hahn spin-echo experiment, consisting of six  $90^\circ$ ,  $\tau$ ,  $180^\circ$  sequences, with  $\tau$  varying from 0.1–1 sec. Note that the amplitude of the echo decreases as  $\tau$  increases. In this experiment the echoes are positive, rather than negative as expected, since a diode detector was used, which measures only amplitude, not phase of the signal. (See Chapter 3 for further details on detectors.)

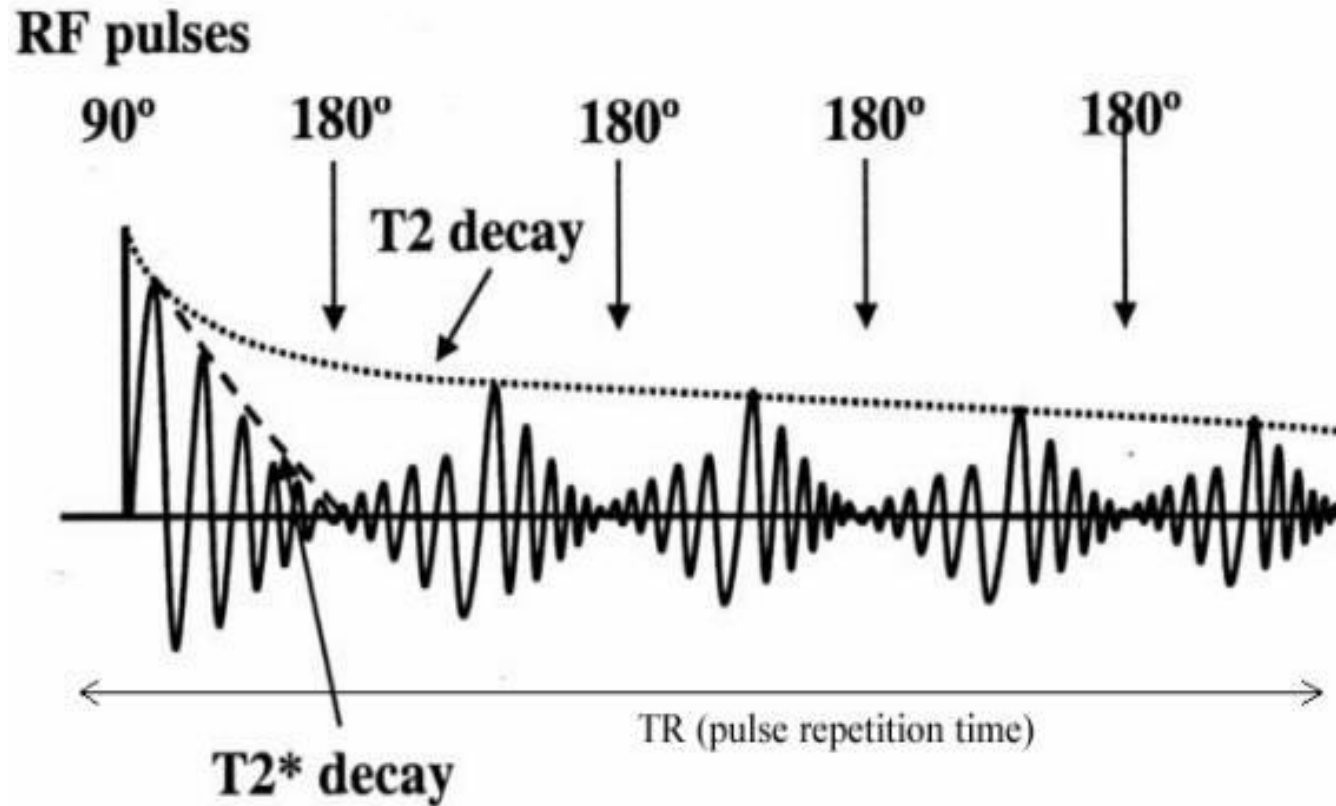
# Carr Purcell sequence



# Carr- Purcell sequence



# Carr-Purcell sequence



- End of the first part

$$[M_0 - M_Z(t)] = [M_0 - M_Z(0)] \cdot e^{-\frac{t}{T_1}}$$

